

A HYBRID COMPUTER FILTER FOR
UNFOLDING GAMMA-RAY SPECTRA

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The author wishes to express his sincere thanks to his major professor, Dr. M. W. Young, for his valuable guidance and counsel given during the course of this work.

A Thesis

during the course of this work.

The author is indebted to Dr. R. C. Martin, for his valuable assistance.
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Agricultural and Mechanical College
in partial fulfillment of the
requirements for the degree of
Master of Science

in

The Department of Nuclear Engineering

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by

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ABSTRACT

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Interpretation of raw gamma spectra from NaI(Tl) scintillation

FIGURE

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empirical scattering function. The unfolding of raw spectra from these required fitting functions was done by the Hyodo model employed by the digital system. Unfolded spectra were obtained by which raw spectra were unfolded in parallel with selection of the peak areas, and then applied to normalization. In the case of Cs-137 spectrum, Unfolded spectra showed experimental features of the peaks, but only small fluctuations, indicating no significant components, just at incident-peaks. The purpose of unfolding was to possibly check the general applicability of the Hyodo model and its validity. Unfolding was done by applying a convolution least-squares method.

ABSTRACT

Interpretation of raw gamma spectra from NaI(Tl) scintillation systems is difficult because of complications introduced by a multiplicity of gamma interactions and light-emission processes within the detector. The application of computers to interpretation permits sophisticated spectral analysis through spectrum stripping or unfolding techniques. Although digital systems have been utilized previously, an analog-digital hybrid system offers the possibility of avoiding laborious and time-consuming parameter selection required for the purely digital approach. The hybrid system was applied to the Hyodo response function procedure by using the analog computer for empirical selection of fifth-order polynomial coefficients for seven required fitting functions; once fitted, the polynomials are employed by the digital system for actual spectrum unfolding. Two calibration spectra were employed for coefficient selection; the procedure was then applied to several "catalog" spectra to illustrate its generality. Unfolded spectra show the expected residual full-energy peaks, with only small Bremsstrahlung, non-rectangular Compton components, and adjacent-peak overlap components as interferences. It is possible that the general approach could be accomplished with a purely digital system by applying analog simulation language.

paper less like a traditional numerical computation device, but more like a "digital computer". The analog computer generates the response parameters for each calculated set of polynomials and communicates this information to the digital computer. These parametric fun-

INTRODUCTION

After an initial effort of potentiometer adjustments, unfolding is per-

Pulse-height distributions obtained from gamma radiation interactions in scintillation spectrometers are complicated representations of monoenergetic, line width, gamma rays (1). Thallium-activated sodium-iodide scintillators have the advantage of high interaction efficiencies in the range of 0 to 5 Mev (approximately 30% at 1 Mev); however, they suffer from spectrum distortion because of limited energy resolution (approximately 7% full-width-half-maximum at 1 Mev). The spectra are further complicated by the statistical nature of the counts in each channel as recorded by a multichannel analyzer.

The several complications, Compton areas, escape peaks, etc., of scintillation gamma-ray spectrometry make the interpretation of complex gamma spectra difficult. Analysis of such spectra has been approached through several computer techniques (2) known as "unfolding", which is an operational process on the measured pulse height distribution to obtain a better approximation of the incident gamma-ray spectrum. Of the various unfolding procedures, a number of them involve analytical and matrix inversion techniques for describing a monoenergetic gamma-ray spectrum with subsequent identification of the troublesome interactions generated within the spectrometer.

The approach presented in this thesis utilizes the digital computer less like a traditional numerical computational device and more like a "digital filter" (3). An analog computer generates the response parametric functions associated with the spectrometer and communicates this information to the filter. These parametric func-

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tions can be readily tuned up for any given experimental set-up with a minimal effort of potentiometer adjustments. Unfolding is performed in the order of milliseconds, and output is in graphical or numerical form, or displayed via digital-to-analog oscilloscope.

rel. to the incident gamma-ray spectrum by the relation:

$$N_{\text{c}} = \frac{N_{\text{t}}}{A_t(E)} \cdot \frac{e_1}{e_2} \cdot \frac{\rho}{\rho_0} \cdot \frac{1}{1 + \frac{E}{E_0}}$$

in which:

N_{t} = The corrected number of counts in each channel of a multichannel counter. The channels correspond to pulse heights E_1, E_2, \dots, E_m .

e_1 = The total initial counting efficiency.

$A_t(E)$ = The trapping efficiency of the detector system, probability that a primary photon is detected at energy E and not scattered with a loss of energy due to absorption along the path.

ρ/ρ_0 = The number of photons in the incident spectrum per unit energy interval dE .

Pulse height distributions may be produced in accordance with the

interactions of the incident gamma-rays with the detector. Pulse height distributions of this type are obtained from the gamma-ray spectrum resulting from the interaction of the primary photon, etc., in the detector and from the signal output of the counter and of the subsequent pulse processing circuit.

It is difficult to make an accurate assessment of a parameter ex-

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without introducing error CHAPTER 1 MODEL (4) outlined resulting spectrum.

The major portion of the gamma scintillation spectrometer is

Triangular Matrix

The pulse-height distribution of a scintillation spectrometer is related to the incident gamma-ray spectrum by the equation:

$$\int_0^{\infty} A_i(E') P(E') dE' = b_i + e_i, \quad i = 1, 2, \dots, n, \quad (1)$$

in which:

b_i = the observed number of counts recorded in the i^{th} channel of a multichannel analyzer, (the channel responds to pulses from V_i to V_{i+1} in height);

e_i = the random statistical counting error;

$A_i(E')$ = the response function of the spectrometer (the probability that a monoenergetic photon source with

energy E' and unit intensity will result in a count in the i^{th} channel); and,

$P(E')dE'$ = the number of photons in the incident spectrum between an energy of E' and $E' + dE'$.

Pulse height distributions are, in truth, a representation of the interactions of the incident gamma flux with the detector. These distributions include distortion and degradation of the true gamma spectrum resulting from Compton scattering, escape peaks, smearing, etc., in the detector itself, and by signal-handling characteristics of the subsequent electronic system.

It is difficult to unfold the response of a spectrometer exactly

without introducing spurious features in the desired resulting spectrum. The major problem of the gamma scintillation spectrometer is not one of instrumental smearing, but rather derives from the fact that the monoenergetic response $A_i(E')$ has spurious peaks resulting from escape processes (pair-production annihilation photons, iodine X-rays, etc.), and a low-energy tail arising from internal and external Compton scattered photons. One approach is to accept the finite instrumental line width (produced by statistical broadening in the crystal and photomultiplier), and then try to remove the tails and spurious escape peaks from the result. The final estimated spectrum obtained will be denoted by $\bar{P}(E)$; this estimated spectrum is related to the true spectrum $P(E')$ by:

$$\bar{P}(E) = \int_0^{\infty} S(E, E') P(E') dE', \quad (2)$$

in which $S(E, E')$ is the inherent smearing function. In other words, $\bar{P}(E)$ is the spectrum with some Gaussian smearing, and $P(E')$ is the true unbroadened spectrum.

The principal problem in obtaining the desired solution, in a discrete formulation, is finding a set of coefficients U_{ik} , $k = 1, 2, \dots, n$, for each of the n response functions $A_i(E')$ such that

$$A_i(E') = \sum_{k=1}^n U_{ik} S(E_k, E'). \quad (3)$$

A significant simplification can be achieved without obtaining numerical values for $A_i(E')$, if one accepts the restriction that a

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linear relation holds for E_k , $k = 1, 2, \dots, n$, to be the energies at which the main peaks of $A_i(E')$ vs E' are centered in channel k . A two dimensional sketch of $A_i(E')$ is given in Figure 1. Note that slices through this response surface parallel to the discrete channel number axis, $A_i(E')$ vs i , are the conventional pulse-height distributions for monoenergetic sources, but that slices through the response surface parallel to the energy axis, $A_i(E')$ vs E' for fixed i , are "efficiency functions" which give the efficiency of a single channel of a gamma photon.

In order to obtain a relation between the observed pulse height distribution, b_i , and the estimated spectrum, $\bar{P}(E)$, a finite set of simultaneous algebraic equations may be derived by substituting Eq. (3) into Eq. (1);

$$\int_0^\infty \left[\sum_{k=1}^n U_{ik} S(E_k, E') \right] P(E') dE' = b_i + e_i. \quad (4)$$

By interchanging the order of integration and summation, the following results:

$$\sum_{k=1}^n U_{ik} \left[\int_0^\infty S(E_k, E') P(E') dE' \right] = b_i + e_i, \quad (5)$$

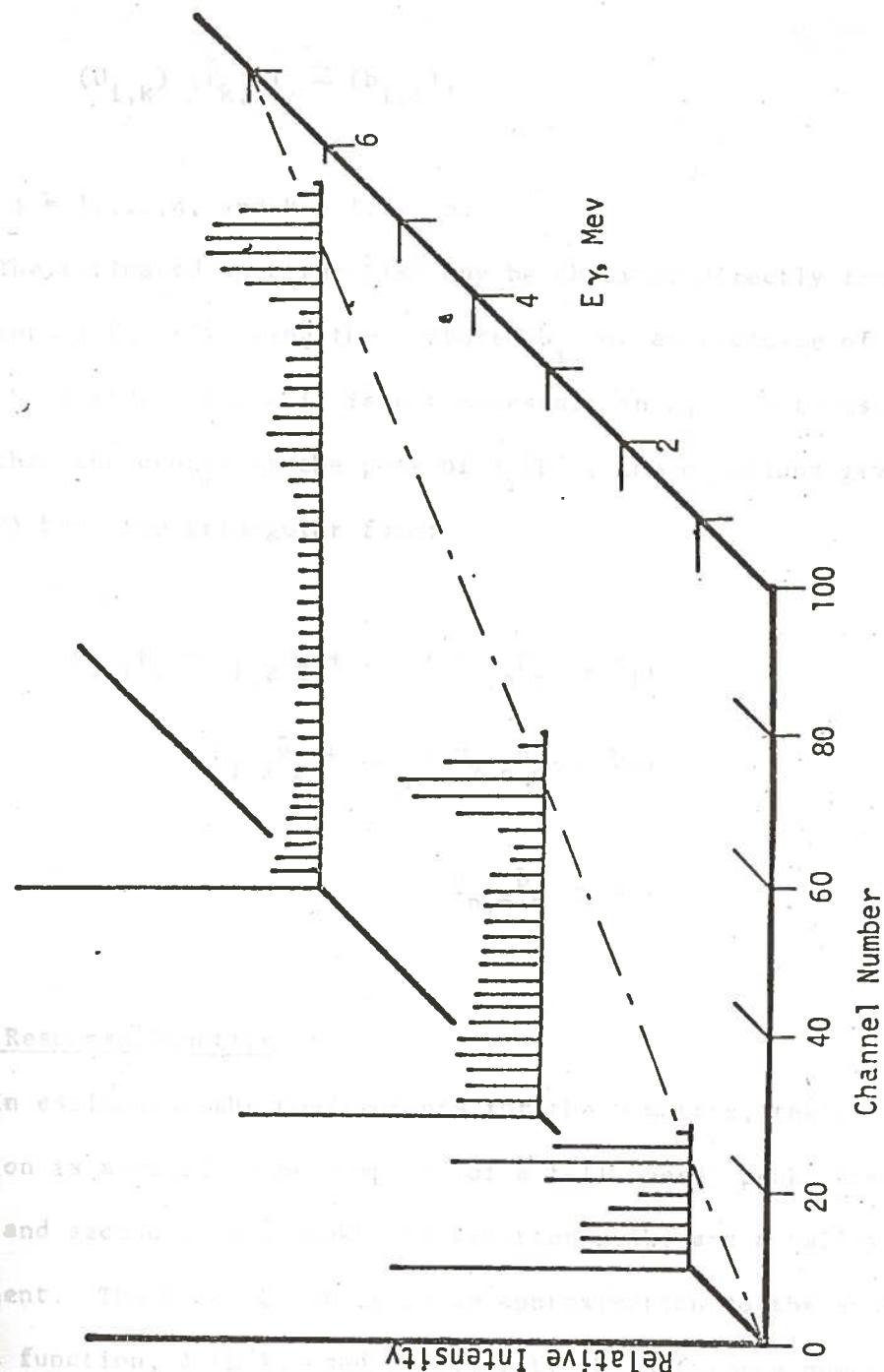
and making use of the definition of $\bar{P}(E)$ given by Eq. (2), the final form is obtained:

$$\sum_{k=1}^n U_{ik} \bar{P}(E_k) = b_i + e_i, \quad i = 1, 2, \dots, n. \quad (6)$$

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Figure 1
Response Function Surface



where $j = 1, \dots, 6$. The ratio of the relative intensities of the peaks is directly proportional to the ratio of the areas of the circles. The ratio of the areas of the circles is given by Eq. (3) for the trapezoidal form.

In this case, the response function is

Hydro Resonance Peaks

In this case, the response function is given by Eq. (3). The first and second derivatives of the response function are zero component. This model results in the following expression for the response function, $R(E)$:

This model results in the following expression for the response function:

The diagonal of this curve has a peak at $E = E_0$ and other peaks at $E = E_1, E_2, \dots, E_6$ from the bottom to the top of the curve.

Finally, expressing Eq. (6) in matrix form and neglecting e_i :

$(U_{i,k}) (\bar{P}_{k,1}) \approx (b_{i,1})$, because the position does not vary much with the energy in the incident photons, as can be seen from the where $i = 1, \dots, n$, and $k = 1, \dots, n$.

The estimated spectrum $\bar{P}(E)$ may be obtained directly from a solution of Eq. (7) using the measured b_i for an estimate of the right hand side. Since it is not necessary in Eq. (7) to use any E_k less than the center of the peak of $A_i(E')$, the equations given by Eq. (7) have the triangular form;

$$\begin{aligned} U_{1,1}\bar{P}_1 + U_{1,2}\bar{P}_2 + \dots + U_{1,n}\bar{P}_n &\approx b_1, \\ U_{2,2}\bar{P}_2 + \dots + U_{2,n}\bar{P}_n &\approx b_2, \\ &\vdots \\ U_{n,n}\bar{P}_n &\approx b_n. \end{aligned}$$

Hyodo Response Function

In estimating the coefficients for the U matrix, the response function is assumed to be composed of a full energy peak, Compton, first and second escape peaks, backscatter peak, and a valley fill component. The Hyodo U matrix is an approximation to the actual response function, $A_i(E')$, used to simplify the unfolding numerically. This model replaces the total energy peaks by isolated components on the diagonal of the matrix. This procedure can also be extended to other peaks in the response function. For instance, the broad escape



peaks are replaced by a sharp ridge of components in the U coefficient matrix. A more complete analysis shows that the backscatter peak must be replaced by several components because its position does not vary much with the energy of the incident photons, as can be seen from the Compton equation for 180° backscatter:

$$E_s = \frac{0.51 E_i}{0.51 + 2E_i},$$

in which E_s is the scattered photon energy and E_i is the incident photon energy. The continuum can be approximated by flat rectangular sections extending up to the pulse height of the Compton edge for the Compton area, and one small rectangle below the total energy peak for a valley fill function. The location of the escape peaks, the backscatter peak and the Compton edge are treated as being functionally related to the photopeak location. These approximations are represented schematically in Figure 2.

The resulting U matrix is composed of five components: (1) Compton fraction, (2) backscatter fraction, (3) first escape fraction, (4) second escape fraction, and (5) valley-fill fraction as follows:

$$\begin{matrix} & U_{1,1} & U_{1,2} & \dots & U_{1,n} \\ & U_{2,2} & \dots & U_{2,n} \\ & & \vdots & & \\ & & & U_{n,n} \end{matrix}$$

where each column is composed of a combination of the components (above) depending upon the position of $U_{diagonal}$. The diagonal is

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1. Compton
 2. First escape peak
 3. Second escape peak
 4. Backscatter peak
 5. Valley-fill function
 6. Full-energy peak

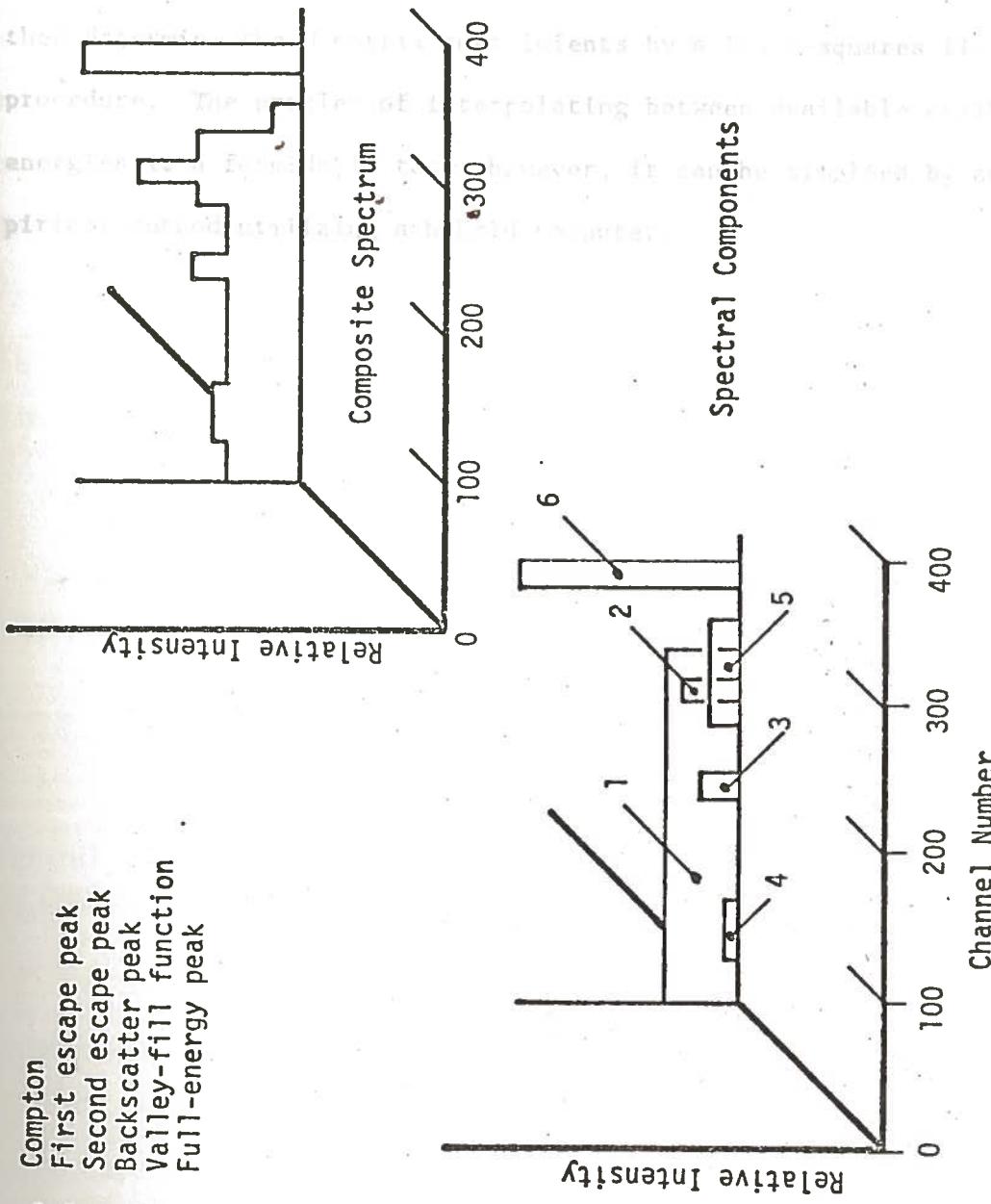


Figure 2
Hyodo Response Model

represented by the peak-to-total ratio.

If many monoenergetic calibration sources are available, one may interpolate by suitable functions between the calibration sources and then determine the U matrix coefficients by a least-squares fitting procedure. The problem of interpolating between available calibration energies is a formidable task, however, it can be simplified by an empirical method utilizing a hybrid computer.

Several programs have been written and are available on the program listing tape. Appendix A contains a listing of the calibration source subprograms.

- (a) Calibration source digital representation routine.
 - (b) Calibration source depth conversion routine.
 - (c) Calibration source energy conversion routine.
 - (d) Calibration source energy conversion routine.
 - (e) Calibration source energy conversion routine.
 - (f) Calibration source energy conversion routine.
 - (g) Calibration source energy conversion routine.
 - (h) Calibration source energy conversion routine.
 - (i) Calibration source energy conversion routine.
 - (j) Calibration source energy conversion routine.
 - (k) Calibration source energy conversion routine.
 - (l) Calibration source energy conversion routine.
 - (m) Calibration source energy conversion routine.
 - (n) Calibration source energy conversion routine.
 - (o) Calibration source energy conversion routine.
 - (p) Calibration source energy conversion routine.
 - (q) Calibration source energy conversion routine.
 - (r) Calibration source energy conversion routine.
 - (s) Calibration source energy conversion routine.
 - (t) Calibration source energy conversion routine.
 - (u) Calibration source energy conversion routine.
 - (v) Calibration source energy conversion routine.
 - (w) Calibration source energy conversion routine.
 - (x) Calibration source energy conversion routine.
 - (y) Calibration source energy conversion routine.
 - (z) Calibration source energy conversion routine.
- Several library subroutines have been written which facilitate the control.

Generating U Matrix Coefficients

To calibrate the system, one must have a minimum of three different sample spectra from the same detector. These spectra are read in and stored in the digital computer. One of the source spectra is then selected for integration. The other two are chosen at random. The black box with variable bandpass filter has a band width of 50 MHz to provide a tuning adjustment. The frequency select feature of the filter

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CHAPTER 2 COMPUTER PROGRAM

A computer program was developed to (1) empirically determine the U matrix coefficients from calibration spectra obtained with a particular spectrometer; (2) unfold pulse-height distributions once the filter is calibrated; (3) output the unfolded spectra via digital plot, numerically, and/or digital-to-analog oscilloscope; and (4) provide overall control between the analog and digital computers. The program (listing in Appendix) consists of a main program plus the following subprograms:

A) Plot - provides digital plot routine, Fortran IV;

B) Scope - provides oscilloscope display of an array stored in the digital computer, Fortran IV and Assembly language;

C) Hybrid- digitizes analog function and stores as an array, Fortran and Assembly language.

Several library (hybrid) routines were utilized for operational control.

Generating U Matrix Coefficients

To calibrate the filter to unfold 4 Mev maximum distributions, sample spectra from the spectrometer (Cs-137,Na-24) are read in and stored in the digital computer. One of the sample spectra is then selected for initial trial. If one considers the hybrid filter as a black box with variable "knobs", the philosophy of design is to provide a tuning adjustment for each objectionable feature of the pulse-

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height distribution which is to be "tuned out". The analog computer is patched (wired) to generate a fifth order polynomial with variable coefficients. Control is shifted to the analog computer and the polynomial is displayed. The five coefficients simulated by potentiometers on the analog computer, can be varied while the resulting Compton function is displayed, and an empirical trial array is digitized and stored in the digital computer. Another five coefficients are obtained for the backscatter function, etc. After all the arrays (Compton, backscatter, first and second escape, valley-fill) are produced, control is transferred to the digital computer for unfolding. The resulting spectrum is filtered by the digitized empirical polynomial function, and displayed on the oscilloscope. Imperfections are then corrected by adjustments of the polynomial coefficients, and the filtering process repeated. After several trials, a nearly perfect filtered spectrum can be displayed. A second calibration spectrum is then selected, and the coefficients for the polynomial are further refined to obtain the spectral components relationship as a function of energy. The ultimate set of coefficients is thus generated by successive trials on several calibration spectra.

Unfolding

Once the response function parameters, in the form of a 5th order polynomial coefficients, are obtained and stored, general unfolding can be initiated for any distribution obtained from the same spectrometer from which the calibration spectra were obtained. Unfolding is accomplished entirely within the digital computer.

After a pulse-height distribution has been read in and stored, and following matrix equation exists:

Response Function

$$\begin{bmatrix} U_{1,1} & U_{1,2} \dots & U_{1,n-1} & U_{1,n} \\ U_{2,2} & \dots & U_{2,n} \\ \vdots \\ U_{n-1,n-1} & U_{n-1,n} \\ U_{n,n} \end{bmatrix} \begin{bmatrix} \bar{P}_1 \\ \bar{P}_2 \\ \vdots \\ \bar{P}_{n-1} \\ \bar{P}_n \end{bmatrix} = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_{n-1} \\ b_n \end{bmatrix}$$

By solving first for \bar{P}_n , and then working backwards toward \bar{P}_1 , the triangular matrix can be reduced to the solution. The results of the unfolding procedure will then be related to the true incident gamma spectrum by Eq. (2). Thus the effect of the Hyodo model is to produce a smeared approximation to the desired spectrum, with the degree of smearing given essentially by the resolution of the detector.

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(potentially) with each of the listed unfolding methods, a hybrid approach.

Using the input listed in Table 37 and the two filters given in Figures 1 and 2, respectively, we present a series of figures which for the Tech Model, we believe, best fit the data. We note that all fulfill the basic requirements of a meaningful physical spectrum except for branching ratios which are unknown. The criteria for setting an incident energy of 100 keV is that the pulse height is 100 times the noise level.

CHAPTER 3 RESULTS

Response Function

Of the nine required functions, the first two (peak-to-total and intrinsic efficiency) were simulated from data obtained from Heath (1) for a 3" x 3" scintillation crystal. The remaining seven functions (Compton, first escape, second escape, backscatter, and three valley-fill functions) were generated as fifth-order polynomials by an empirical selection of potentiometer settings for the five necessary coefficients. Cesium-137 and Na-24 spectra were employed for progressive trial and error refinements of the necessary potentiometer settings. These spectra were chosen because they provide a wide energy range and relatively uncomplicated distributions.

The general shapes of the seven empirically fitted polynomials are presented in Figures 3 and 4. The functional coefficients (potentiometer settings) can be selected within 30 minutes utilizing a hybrid computer.

Unfiltered and filtered (Cs-137 and Na-24) spectra are illustrated in Figures 5 and 6, respectively, to present a measure of "goodness" for the technique. As apparent from the figures, the unfolded spectra fulfill the objectives of removal of the complicating features, except for Bremsstrahlung in the Na-24 distribution. The criteria for setting an individual channel to zero was plus or minus the reciprocal of the plotting scale height (e.g., = 0.01 for a scale height of 100) times the maximum count.

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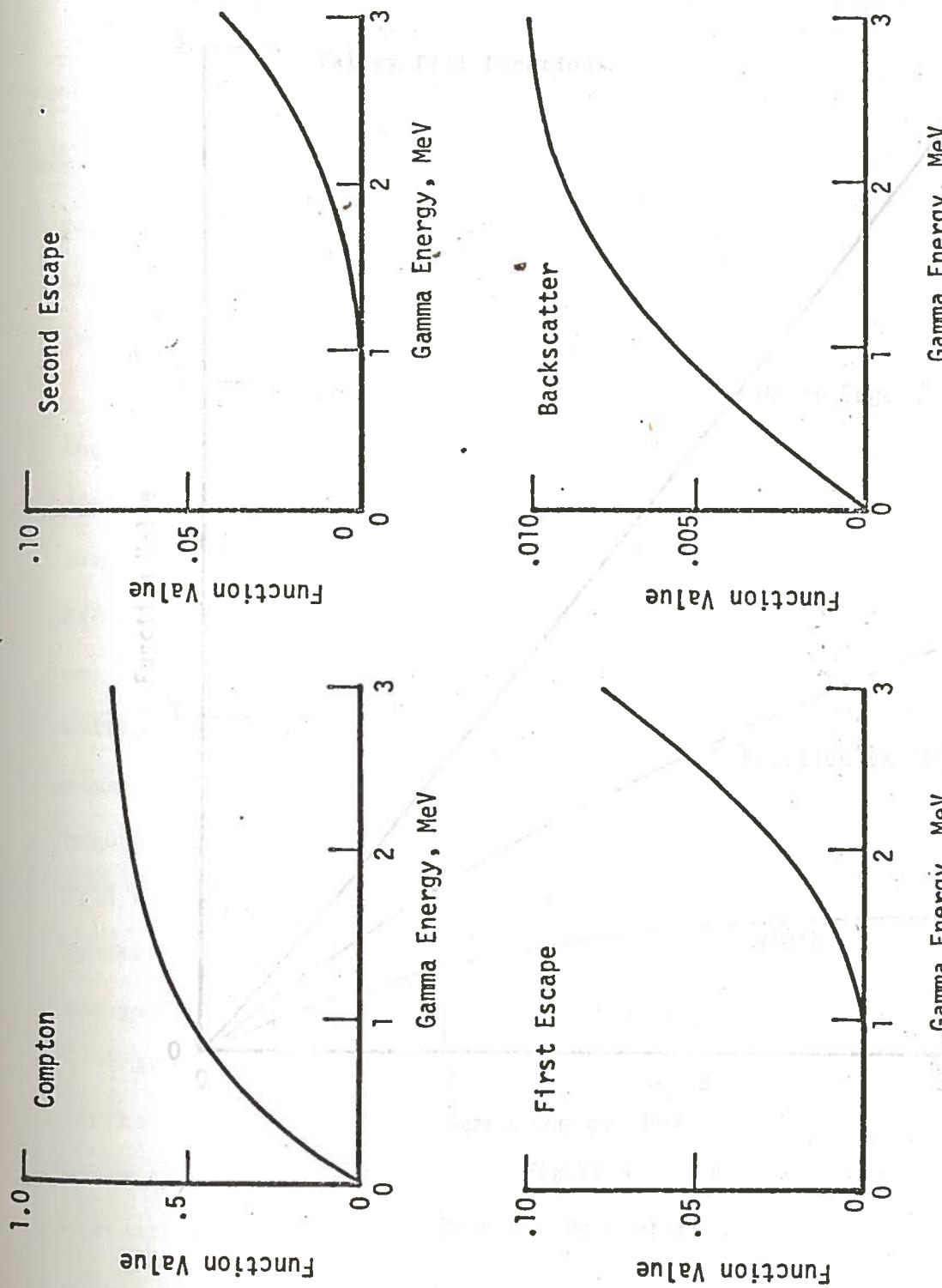


Figure 3
Response Parameters

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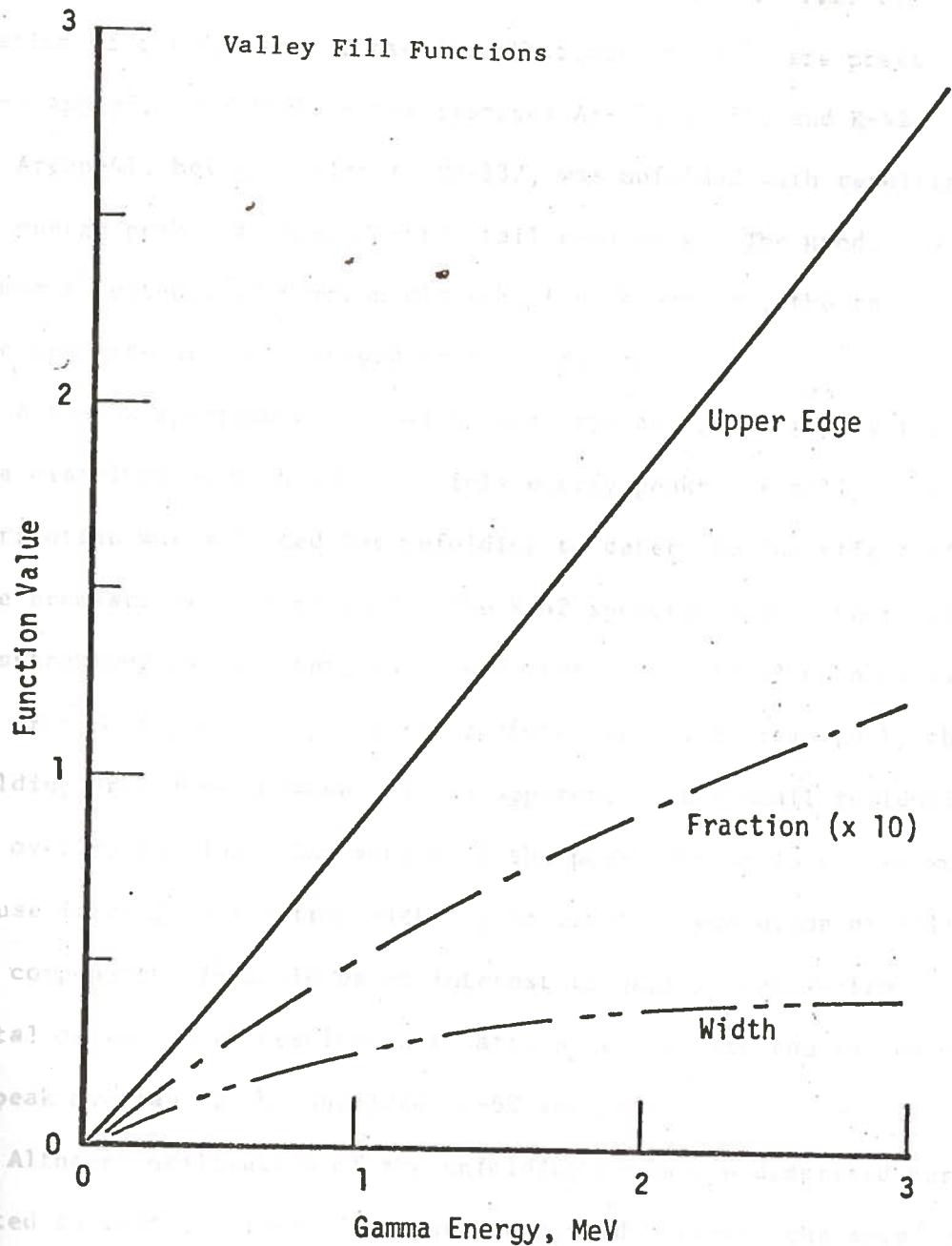


Figure 4

Response Parameters

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Filter

Additional spectra from Beach (1) were unfolded to test the operation of the filter. These distributions (folded) are presented in the Appendix and include the isotopes Ar-41, Co-60, and K-42.

Argon-41, being similar to Cs-137, was unfolded with resulting full energy peak and upper Compton tail remaining. The Hyodo model assumes a rectangular Compton distribution; therefore, the true Compton components are not removed from the spectrum.

A Co-60 spectrum was tried to test the applicability of filtering a distribution with adjacent full energy peaks. Finally, a K-42 distribution was selected for unfolding to determine the effect of a large Bremsstrahlung component. The K-42 spectrum shows the residual Bremsstrahlung as expected, and demonstrates the effect of a continuum. For Co-60, the complicating features were well removed by the unfolding procedure; however, it is apparent that a small residual peak overlap remains. The source of the peak overlap is not known because it could arise from either poor crystal resolution or valley fill component. It would be of interest to employ spectra from a crystal of very high resolution to attempt to isolate the source of the peak overlap of the unfolded Co-60 spectrum.

Although utilization of the unfolding procedure described here is limited to installations of comparable hybrid systems, the speed with which the response functions can be fitted makes this operation attractive. It is possible that the procedure could be applied to purely digital systems by introduction of an analog simulation language.

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Figure 5

Unfolded Cs-137 Spectrum
Pages 19 and 20

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CHAN SPECT

CS-137

	0.	.
1	0.	0
2	1483.	.
3	25821.	.
4	39892.	.
5	1746.	.
6	0.	0
7	0.	0
8	0.	0
9	0.	0
10	0.	0
11	0.	0
12	0.	0
13	0.	0
14	0.	0
15	0.	0
16	0.	0
17	0.	0
18	0.	0
19	0.	0
20	0.	0
21	0.	0
22	0.	0
23	0.	0
24	0.	0
25	0.	0
26	0.	0
27	0.	0
28	0.	0
29	0.	0
30	0.	0
31	0.	0
32	0.	0
33	0.	0
34	0.	0
35	0.	0
36	0.	0
37	0.	0
38	0.	0
39	0.	0
40	0.	0
41	0.	0
42	0.	0
43	0.	0
44	0.	0
45	0.	0
46	0.	0
47	0.	0
48	0.	0
49	0.	0
50	0.	0
51	0.	0
52	0.	0

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55 0. 0
56 0. 0
57 0. 0
58 0. 0
59 1005. 0
60 3201. .*
61 6077. . *
62 11844. . *
63 22252. .
64 35252. .
65 46975. .
66 53404. .
67 49923. .
68 39053. .
69 25915. .
70 14483. . *
71 6973. . *
72 2942. .*
73 1210. .
74 0. 0
75 0. 0
76 0. 0
77 0. 0
78 0. 0
79 0. 0
80 0. 0
81 0. 0
82 0. 0
83 0. 0
84 0. 0
85 0. 0
86 0. 0
87 0. 0
88 0. 0
89 0. 0
90 0. 0
91 0. 0
92 0. 0
93 0. 0
94 0. 0
95 0. 0
96 0. 0
97 0. 0
98 0. 0
99 0. 0
100 0. 0

STOP 0

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Figure 5a

Folded Cs-137 Spectrum
Pages 22 and 23

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CHAN SPECT

CS-137

	0	.	*
1	6042.	.	*
2	7506.	.	*
3	31921.	.	*
4	45993.	.	*
5	7778.	.	*
6	6031.	.	*
7	6009.	.	*
8	5955.	.	*
9	5645.	.	*
10	5721.	.	*
11	5849.	.	*
12	5750.	.	*
13	5694.	.	*
14	5933.	.	*
15	5950.	.	*
16	5920.	.	*
17	6509.	.	*
18	6884.	.	*
19	7071.	.	*
20	6773.	.	*
21	6499.	.	*
22	6205.	.	*
23	6081.	.	*
24	6096.	.	*
25	6267.	.	*
26	6247.	.	*
27	6234.	.	*
28	6124.	.	*
29	6103.	.	*
30	6086.	.	*
31	6010.	.	*
32	6003.	.	*
33	5879.	.	*
34	5951.	.	*
35	5904.	.	*
36	5935.	.	*
37	5871.	.	*
38	5977.	.	*
39	5947.	.	*
40	5979.	.	*
41	6172.	.	*
42	6271.	.	*
43	6534.	.	*
44	6416.	.	*
45	6129.	.	*
46	5893.	.	*
47	4837.	.	*
48	3877.	.	*
49	2959.	.	*
50	2330.	.	*
51	1957.	.	*
52	1622.	.	*

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55	1219.	.
56	1301.	.
57	1287.	.
58	1498.	.
59	1933.	.
60	3201.	..*
61	6077.	.
62	11844.	.
63	22252.	.
64	35252.	.
65	46976.	.
66	53404.	.
67	49923.	.
68	39055.	.
69	25915.	.
70	14483.	.
71	6979.	.
72	2942.	..*
73	1210.	.
74	451.	0
75	211.	0
76	89.	0
77	32.	0
78	63.	0
79	38.	0
80	0.	0
81	0.	0
82	0.	0
83	0.	0
84	0.	0
85	0.	0
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87	0.	0
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97	0.	0
98	0.	0
99	0.	0
100	0.	0

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CHART SHEET



Figure 6

Unfolded Na-24 Spectrum

Pages 25 through 30

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CHAN SPECT

	0
1	0.	0
2	5421.	*
3	4024.	*
4	2173.	*
5	1104.	*
6	521.	0
7	0.	0
8	0.	0
9	0.	0
10	0.	0
11	0.	0
12	0.	0
13	0.	0
14	0.	0
15	0.	0
16	0.	0
17	0.	0
18	0.	0
19	0.	0
20	0.	0
21	0.	0
22	0.	0
23	0.	0
24	0.	0
25	0.	0
26	0.	0
27	0.	0
28	0.	0
29	0.	0
30	0.	0
31	0.	0
32	0.	0
33	0.	0
34	0.	0
35	0.	0
36	0.	0
37	0.	0
38	0.	0
39	0.	0
40	0.	0
41	0.	0
42	0.	0
43	0.	0
44	0.	0
45	0.	0
46	0.	0
47	0.	0
48	0.	0
49	0.	0
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51	0.	0
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NA-24

NUCLEAR SCIENCE CENTER
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NS

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57 0. 0
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124	1732. .*
125	2352. .*
126	4001. .
127	6675. .
128	9233. .
129	13451. .
130	17937. .
131	23765. .
132	28817. .
133	30642. .
134	30489. .
135	28392. .
136	24328. .
137	20082. .
138	15035. .
139	10719. .
140	7305. .
141	4553. .
142	2535. .
143	1329. .*
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NUCLEAR SCIENCE CENTER
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JOHNSON SCIENCE CENTER
LOUISIANA STATE UNIVERSITY

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246	0.	0
247	0.	0
248	0.	0
249	0.	0
250	0.	0
251	321.	0
252	1002.	.
253	2216.	• *
254	3013.	• *
255	3679.	• *
256	4516.	• *
257	5483.	• *
258	6593.	• *
259	7715.	• *
260	8693.	• *
261	9673.	• *
262	10165.	• *
263	10653.	• *
264	10392.	• *
265	10248.	• *
266	9736.	• *
267	8953.	• *
268	7832.	• *
269	6852.	• *
270	5882.	• *
271	5024.	• *
272	4225.	• *
273	3366.	• *
274	2662.	• *
275	2122.	• *
276	1685.	• *

279 782. .
280 0. 0
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Figure 6a

Folded Na-24 Spectrum

Pages 32 through 37

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CHAN SPECT

CHAN	SPECT
1	8969.
2	13493.
3	12200.
4	10384.
5	9339.
6	8793.
7	8337.
8	8207.
9	7998.
10	7799.
11	7510.
12	7463.
13	7364.
14	7279.
15	7200.
16	7227.
17	7104.
18	7285.
19	7531.
20	7613.
21	8332.
22	8712.
23	9103.
24	8702.
25	8431.
26	7650.
27	7407.
28	7146.
29	7044.
30	7003.
31	6983.
32	7019.
33	7067.
34	7026.
35	7055.
36	6983.
37	6791.
38	6613.
39	6521.
40	6497.
41	6497.
42	6417.
43	6387.
44	6472.
45	6343.
46	6410.
47	6511.
48	6614.
49	6953.
50	7347.
51	7534.
52	7400.

32
33

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CMI

55 6353. • *
56 6233. • *
57 6219. • *
58 6089. • *
59 6136. • *
60 6162. • *
61 5970. • *
62 5999. • *
63 6144. • *
64 6169. • *
65 6226. • *
66 6153. • *
67 6220. • *
68 6203. • *
69 6165. • *
70 6219. • *
71 6249. • *
72 6011. • *
73 6343. • *
74 6443. • *
75 6433. • *
76 6463. • *
77 6591. • *
78 6475. • *
79 6615. • *
80 6690. • *
81 6783. • *
82 6843. • *
83 7142. • *
84 7266. • *
85 7403. • *
86 7353. • *
87 7483. • *
88 7481. • *
89 7410. • *
90 7383. • *
91 7522. • *
92 7609. • *
93 7745. • *
94 7697. • *
95 7646. • *
96 7683. • *
97 7697. • *
98 8017. • *
99 8005. • *
100 8032. • *
101 8115. • *
102 8233. • *
103 8161. • *
104 8227. • *
105 8273. • *
106 8404. • *
107 8175. • *
108 8306. • *

LOUISIANA STATE UNIVERSITY
LIBRARIES

CAI

111 7323. • *
112 6653. • *
113 6084. • *
114 5444. • *
115 4967. • *
116 4479. • *
117 3962. • *
118 3634. • *
119 3444. • *
120 3337. • *
121 3252. • *
122 3226. • *
123 3462. • *
124 3767. • *
125 4620. • *
126 6111. • *
127 8173. • *
128 11513. • *
129 15896. • *
130 20255. • *
131 25953. • *
132 30283. • *
133 32949. • *
134 32855. • *
135 30693. • *
136 26741. • *
137 22299. • *
138 17296. • *
139 13085. • *
140 9573. • *
141 6824. • *
142 4886. • *
143 3803. • *
144 2985. • *
145 2585. • *
146 2349. • *
147 2239. • *
148 2153. • *
149 2016. • *
150 2100. • *
151 2121. • *
152 1963. • *
153 2057. • *
154 2049. • *
155 2071. • *
156 2141. • *
157 2182. • *
158 2271. • *
159 2280. • *
160 2397. • *
161 2631. • *
162 2941. • *
163 3093. • *
164 3391. • *

167 3745. . *
168 3741. . *
169 3713. . *
170 3682. . *
171 3526. . *
172 3688. . *
173 3413. . *
174 3253. . *
175 3211. . *
176 2992. . *
177 3055. . *
178 3043. . *
179 2975. . *
180 2882. . *
181 3025. . *
182 2952. . *
183 2939. . *
184 3022. . *
185 3000. . *
186 3074. . *
187 3028. . *
188 3104. . *
189 3059. . *
190 3147. . *
191 3063. . *
192 3117. . *
193 3141. . *
194 3114. . *
195 3207. . *
196 3145. . *
197 3183. . *
198 3105. . *
199 3138. . *
200 3146. . *
201 3214. . *
202 3072. . *
203 3122. . *
204 3116. . *
205 3209. . *
206 3263. . *
207 3301. . *
208 3524. . *
209 3623. . *
210 4142. . *
211 4339. . *
212 4737. . *
213 5307. . *
214 5517. . *
215 5903. . *
216 6247. . *
217 6462. . *
218 6585. . *
219 6433. . *
220 6202. . *

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CMI

223 5298. *
224 5153. *
225 5103. *
226 4700. *
227 4602. *
228 4437. *
229 4253. *
230 4187. *
231 4150. *
232 4111. *
233 4126. *
234 4067. *
235 4128. *
236 3933. *
237 3763. *
238 3767. *
239 3573. *
240 3373. *
241 3252. *
242 2969. *
243 2867. *
244 2641. *
245 2435. *
246 2225. *
247 2056. *
248 1993. *
249 1864. *
250 1923. *
251 1973. *
252 2167. *
253 2515. *
254 3013. *
255 3673. *
256 4516. *
257 5483. *
258 6593. *
259 7715. *
260 8693. *
261 9679. *
262 10165. *
263 10653. *
264 10392. *
265 10243. *
266 9736. *
267 8953. *
268 7832. *
269 6852. *
270 5382. *
271 5024. *
272 4225. *
273 3366. *
274 2662. *
275 2123. *
276 1685. *

279 782. 0
280 580. 0
281 443. 0
282 344. 0
283 267. 0
284 260. 0
285 205. RO 1.0. Rehabilitation Project, 1972, Research and
286 182. 0 (Journal of Rehabilitation, 1974, 31st Edition).
287 163. 0 (Journal, August 1974).
288 142. 0
289 130. 0 (Journal, October 1974, Special Edition for
290 154. 0 Rehabilitation Project, 1972, Research and
291 123. 0 (Journal, November 1974). (See also 285, 1972),
292 97. 0
293 97. 0
294 99. 0 (Journal, December 1974, Special Edition for
295 105. 0 Rehabilitation Project, 1972, Research and
296 93. 0 (Journal, January 1975).
297 86. 0
298 86. 0 (Journal, February 1975). (See also 285, 1972)
299 99. 0 (Journal, March 1975). (See also 285, 1972)
300 103. 0 (Journal, April 1975). (See also 285, 1972)

STOP C

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3. Young, M. H. and Burrus, W. R., "A Digital Filter for Unfolding Pulse-Height Distributions", Nuclear Instruments and Methods, 62 (1968), 82-92.
4. Burrus, W. R. and Young, M. H., "Sego: A Computer Code for Unfolding Experimental Pulse-Height Distributions", AEC Research and Development Report, ORNL-TM-2172, July 1968.

Louisiana State University

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Louisiana State University

APPENDIX

TEXAS STATE UNIVERSITY

41

APPENDIX A
COMPUTER PROGRAM LISTING
Pages 41 through 52

```

DIMENSION X(400),ISPECT(400),ID(5)
DIMENSION EFF(400),PT(400),COM(400),BSP(400),
1VFP(400),VFPH(400),VFED(400),FESC(400),SESC(400)
DIMENSION E(6),P(6),C(6),R(6),
1VP(6),VW(6),VE(6),F(6),S(6)
DIMENSION IPOLY(4),POLYVL(4),IRES(6)
LOGICAL TRSL
DATA IPOLY/.64,.60,.65,.70/
DATA POLYVL/.25,.50,.75,.125/
DATA IRES/62,67,72,77,82,100/
K=0
NRESP=400
D8500I=1,4
C      SETS UP POLYNOMIAL
C
500 CALL SPST(IPOLY(I),POLYVL(I))
C      READ IN POLYNOMIAL COEFFICIENTS
C
READ(5,508)(E(I),I=1,6)
READ(5,508)(P(I),I=1,6)
READ(5,508)(C(I),I=1,6)
READ(5,508)(R(I),I=1,6)
READ(5,508)(VP(I),I=1,6)
READ(5,508)(VE(I),I=1,6)
READ(5,508)(VW(I),I=1,6)
READ(5,508)(F(I),I=1,6)
READ(5,508)(S(I),I=1,6)
508 FORMAT(6F10.4)
C      READ IN SPECTRUM ID AND CALIBRATION
C
100 READ(5,105)IL,NCHAN,CHANH
105 FORMAT(5A4,I10,F10.5)
IF(NCHAN.EQ.0)GOTO999
C      READ IN SPECTRUM
C
READ(5,110)(ISPECT(I),I=1,NCHAN)

```

110 FORMAT(12X,10I6)
C
C WRITE SPECTRUM
C
C WRITE(6,111)ID
111 FORMAT(J1',35X,5A4,//)
WRITE(6,112)(ISPECT(I),I=1,NCHAN)
112 FORMAT(10I12)
D0120I=1,NCHAN
120 X(I)=ISPECT(I)
C
C DIGITAL PLOT OF INPUT SPECTRUM
C
CALL PLOT(NCHAN,1D,X)
CALL SAM9('1C')
WRITE(102,125)
125 FORMAT(//)
READ(101,126)
126 FORMAT(I1)
C
C DISPLAY INPUT SPECTRUM VIA OSCILLOSCOPE
C
CALL SCOPE(NCHAN,X)
IF(K.EQ.2)GBT\$600
C
C RESPONSE FUNCTION ROUTINE
C
C
510 WRITE(6,511)
511 FORMAT('1EFF',//)
CALL SSRM(64)
CALL HYBRID(K,NRESP,IRFS,E,ISPECT,EFF)
D0513I=1,NRESP
513 EFF(I)=EFF(I)/20.
WRITE(6,509)(EFF(I),I=1,NRESP)
509 FORMAT(10F10.4)
515 WRITE(6,516)
516 FORMAT('1PT',//)
CALL HYBRID(K,NRESP,IRFS,P,ISPECT,PT)

WRITE(6,509)(PT(I),I=1,NRESP)
IF(K.EQ.1)G0T0580
520 WRITE(6,521)
521 FORMAT('1C9M',//)
CALL SSRP(64)
CALL HYBRID(K,NRESP,IRES,C,ISPECT,C9M)
WRITE(6,509)(C9M(I),I=1,NRESP)
IF(K.EQ.1)G0T0580
525 WRITE(6,526)
526 FORMAT('1BSP',//)
CALL SSRP(64)
CALL HYBRID(K,NRESP,IRES,B,ISPECT,BSP)
D9527I=1,NRESP
527 BSP(I)=BSP(I)/10.
WRITE(6,509)(BSP(I),I=1,NRESP)
IF(K.EQ.1)G0T0580
530 WRITE(6,531)
531 FORMAT('1VFP',//)
CALL SSRP(64)
CALL HYBRID (K,NRESP,IFES,VP,ISPECT,VFP)
WRITE(6,509)(VFP(I),I=1,NRESP)
535 WRITE(6,536)
536 FORMAT('1VFED',//)
CALL HYBRID (K,NRESP,IRES,VE,ISPECT,VFED)
D9538I=1,NRESP
538 VFED(I)=VFED(I)*10.
WRITE(6,509)(VFED(I),I=1,NRESP)
540 WRITE(6,541)
541 FORMAT('1VFWH',//)
CALL HYBRID (K,NRESP,IRES,VW,ISPECT,VFWH)
WRITE(6,509)(VFWH(I),I=1,NRESP)
IF(K.EQ.1)G0T0580
545 WRITE(6,546)
546 FORMAT('1FESC',//)
CALL SSRP(64)
CALL HYBRID (K,NRESP,IRES,F,ISPECT,FESC)
WRITE(6,509)(FESC(I),I=1,NRESP)
550 WRITE(6,551)
551 FORMAT('1SESC',//)
CALL HYBRID (K,NRESP,IRES,S,ISPECT,SESC)

```
      WRITE(6,509)(SESC(I),I=1,NRESP)
580  CALL SAMR('IC1')
      D9900I=1,NCHAN
590  ISPECT(I)=X(I)
C
C      UNFOLDING ROUTINE
C
C      DIAGONAL CALCULATION
C
800  D9900I=1,NCHAN
      KT=NCHAN-I
      CHAN=KT
C
C      ESCAPE PEAK CALCULATION
C
      ESC=.511/CHANH
      KESC1=CHAN-ESC
      KESC2=CHAN-23*ESC
C
C      COMPTON EDGE CALCULATION
C
      ENERGY=CHAN*CHANH
      EDGE=ENERGY/(1.+.2555/ENERGY)
      EDGECH=EDGE/CHANH
      KEDCH=EDGECH
      COMHT=COM(KT)/EDGECH
C
C      VALLEY FILL CALCULATION
C
      IFILED=VFED(KT)/CHANH
      IFILRH=VFWH(KT)/CHANH
      IFILBE=IFILED-IFILRH
      VFNC=VFWH(KT)/CHANH
      VFHT=VFP(KT)/VFNC
C
C      BACKSCATTER CALCULATION
C
      KB=KT-KEDCH
```

KBEND=KB+3
KBSE=KB-3
C
C NORMALIZATION WITH PEAK-TO-TOTAL RATIO
C
XPT=X(KT)/PT(KT)
IF(KT-1)900,900,305
C
C COMPTON FILTER
C
805 IF(KEDCH)850,850,810
810 D9315L=1,KEDCH
815 X(L)=X(L)-XPT*CSHT
C
C ESCAPE PEAKS FILTER
C
820 IF(KESC1)840,840,820
X(L)=X(L)-XPT*FESC(KT)
IF(KESC2)840,840,825
825 X(L)=X(L)-XPT*GESC(KT)
C
C BACKSCATTER FILTER
C
840 IF(KS)850,850,845
845 D9346L=KBSE,KBEND
846 X(L)=X(L)-XPT*BSP(KT)
C
C VALLEY FILL FILTER
C
250 IF(IFILED)900,900,955
855 D9260L=IFILBE,IFILED
860 X(L)=X(L)-XPT*VFHT
900 CONTINUE
CALL SCRE(NCHAN,X)
C
C CONTROL ROUTINE
C
C
IF(TRSL(0)+AND+TRSL(1))GAT9906

```
906 IF(TRSL(0)•AND•TRSL(1))GOT9900
      K=1
      GOT07I=1,NCHAN
907 X(I)=ISPECT(I)
      IF(TRSL(0))GOT9910
910 IF(TRSL(0))GOT9510
      IF(TRSL(1))GOT9911
911 IF(TRSL(1))GOT9520
      IF(TRSL(2))GOT9912
912 IF(TRSL(2))GOT9525
      IF(TRSL(3))GOT9913
913 IF(TRSL(3))GOT9530
      IF(TRSL(4))GOT9914
914 IF(TRSL(4))GOT9545
      GOT0800
990 CALL PLOT(NCHAN, ID, X)
      K=2
      GOT0100
999 CALL RELECE
      END
```

```
SUBROUTINE PLGT(NCHAN,TD,Z)
DIMENSION Z(400)
DIMENSION Y(111)
DIMENSION ID(5)
DATA BLANK,DET,STAR,ZERO/''*','*','*','*','*','*'/
700 CMAX=0.0
D9705I=1,NCHAN
TERM=Z(I)
IF(TERM.GT.CMAX)CMAX=TERM
705 CONTINUE
WRITE(6,710)ID
710 FORMAT('1CHAN',1X,'SPECT',1X,'COUNT',20X,5A4)
D9715I=1,111
715 Y(I)=DET
Y(11)=ZERO
WRITE(6,720)(Y(J),J=1,111)
720 FORMAT(19X,111A1)
D9730I=1,111
730 Y(I)=BLANK
D9750I=1,NCHAN
L=(Z(I)*100.)/(CTMAX)
Y(L+11)=STAR
Y(11)=DET
IF(L.EQ.0)Y(11)=ZERO
WRITE(6,740)I,Z(I),L,(Y(J),J=1,111)
740 FORMAT(I4,1X,F5.3,1X,I4,3X,111A1)
750 Y(L+11)=BLANK
RETURN
END
```

SUBROUTINE SCOPE(NCHAN,Z)
DIMENSION Z(400)
LOGICAL TRSL
DISMAX=0.0
DB401 I=1,NCHAN
ANS=Z(I)
IF(ANS.GT.DISMAX)DISMAX=ANS
401 CONTINUE
400 DB405 I=1,NCHAN
FRACT=Z(I)/DISMAX
IDANS=FRACT+32767
S WD,0 X'E131'
S LW,1 I
S LW,2 *NCHAN
S CI,1 1
S BNE 810
S LI,10 X'0400'
S WD,10 X'E121'
S10 LCW,5 IDANS
S WD,5 X'E065'
S CW,2 1
S BNE 8405
S LI,10 X'0400'
S WD,10 X'E123'
405 CONTINUE
IF(TRSL(5))GOT8410
410 IF(.NOT.TRSL(5))GOT8400
RETURN
END

```
SUBROUTINE HYBRID(K,NRESP,IRES,A,ISPECT,Z)
DIMENSION IRES(6),A(6),ISPECT(400),Z(400)
LOGICAL TRSL
K1=0
300 D9310I=1,6
310 CALL SPOT(IRES(I),A(I))
IF(K.EQ.1)GOT9328
312 CALL PARA
CALL HYAIT(S,NRESP,ISPECT,KH)
315 IF(KH.LT.NRESP)GOT9315
CALL SAMO('IC')
D9320I=1,NRESP
320 Z(I)=FLBAT(ISPECT(I))/32767.
IF(K.EQ.0)RETURN
IF(K1.EQ.5)RETURN
328 CALL STCB('NM')
CALL SAMO('PF')
IF(TRSL(5))GOT9330
330 IF(.NOT.TRSL(5))GOT9330
CALL SAMB('PC')
D9335I=1,6
CALL SACB('PI',IRES(I))
335 CALL DVMR(A(I))
WRITE(6,325)(IRES(I),A(I),I=1,6)
325 FORMAT(110,F10.4)
K1=5
G9TE312
END.
```

	DEF	HYAIT
HYAIT	WD,3	X'0037'
	LI,3	X'FE00'
	WD,3	X'11102'
	LW,3	XPSD
	XN,3	-X'60'
	STH,3	SAVE
	LW,3	*15
	LW,3	*3
	WD,3	X'E134'
	MTW,1	15
	LW,3	*15
	BCR,1	\$+2
	LW,3	*3
	LW,3	*3
	STA,3	COUNT
	MTW,1	15
	LW,3	*15
	BCR,1	\$+2
	LW,3	*3
	AND,3	WORD
	BR,3	STS
	STR,3	STORE
	MTW,1	15
	LW,3	*15
	STW,3	ADDR
	LI,3	0
	STW,3	*ADDR
G8	RD,3	X'E226'
	CW,3	MASK
	BANZ	G8
	MTW,1	15
	LI,3	X'8000'
	WD,3	X'1202'
	WD,2	X'0021'
	LW,3	WNE
	WD,3	X'E121'
	B	*15
XPSD	GEN,8,7,17	X'CF',0,TRANS
	BOUND	8

TRANS	DATA	C,0,INT,0
STD	DATA	X'35300000'
WORD	DATA	X'0001FFFF'
COUNT	DATA	0
MASK	DATA	X'0004'
SAVE	DATA	0
REG	DATA	0
ADDR	DATA	0
WONE	DATA	X'0001'
INT	STW,3	REG
	WD,3	X'E135'
	MTW,1	*ADDR
TEST	RD,3	X'E226'
	CD,3	MASK
	BANZ	TEST
	RD,3	X'E235'
STARE	DATA	C
	MTW,1	STARE
	MTW,-1	COUNT
	BE	XIT
RETURN	LW,3	REG
	GEN,12,3,17	X'0E31,0,TRANS
XIT	LI,3	X'7E00'
	WD,3	X'1302'
	LW,3	WONE
	WD,3	X'E123'
	LW,3	SAVE
	STW,3	X'60'
	WD,3	X'0027'
	LW,3	INHBT
	AND,3	TRANS+1
	STW,3	TRANS+1
	B	RETURN
INHBT	DATA	X'F8FFFFFF'
	END	

	DEF	PARA
PARA	LI,6	X'0300'
	WD,6	X'E02A'
	LI,6	0
	WD,6	X'E02B'
	WD,6	X'E02B'
	LI,6	X'02C0'
	WD,6	X'E02B'
	LI,6	X'0100'
	WD,6	X'E02A'
	B	415
	END	

APPENDIX B.**SELECTED PULSE-HEIGHT DISTRIBUTIONS****Pages 54 through 66**

CHAN SPECT

A-41

	0.....
1	12000.
2	12735.
3	15079.
4	14793.
5	16286.
6	14047.
7	12447.
8	11130.
9	10432.
10	9673.
11	9274.
12	9093.
13	9080.
14	8782.
15	8802.
16	8933.
17	8946.
18	9266.
19	9547.
20	10457.
21	11295.
22	11077.
23	10329.
24	9442.
25	9013.
26	8942.
27	8785.
28	8598.
29	8375.
30	8314.
31	8213.
32	8190.
33	8161.
34	8133.
35	8085.
36	8193.
37	8011.
38	7968.
39	7969.
40	8193.
41	8150.
42	7911.
43	8013.
44	8011.
45	8053.
46	8048.
47	7959.
48	7917.
49	7900.
50	8027.
51	8105.
52	8143.

55	7471.	*
56	8005.	*
57	8341.	*
58	8010.	*
59	8087.	*
60	8240.	*
61	8278.	*
62	8270.	*
63	8335.	*
64	8523.	*
65	8516.	*
66	8444.	*
67	8414.	*
68	8540.	*
69	8684.	*
70	8538.	*
71	8725.	*
72	8870.	*
73	8813.	*
74	9205.	*
75	9277.	*
76	9351.	*
77	9647.	*
78	9685.	*
79	9879.	*
80	9903.	*
81	9985.	*
82	9933.	*
83	10002.	*
84	10059.	*
85	10361.	*
86	10357.	*
87	10621.	*
88	10432.	*
89	10651.	*
90	10814.	*
91	10829.	*
92	10985.	*
93	11185.	*
94	11232.	*
95	11523.	*
96	11526.	*
97	11932.	*
98	11962.	*
99	11833.	*
100	11743.	*
101	11489.	*
102	10894.	*
103	10006.	*
104	9138.	*
105	8019.	*
106	6972.	*
107	6031.	*
108	5024.	*

111 3030. .
112 2932. .
113 2543. .
114 2456. .
115 2603. .
116 3005. .
117 3310. .
118 5127. .
119 8212. .
120 12630. .
121 19207. .
122 28111. .
123 38032. .
124 47623. .
125 54143. .
126 56900. .
127 55161. .
128 49739. .
129 41269. .
130 31509. .
131 22749. .
132 15573. .
133 10108. .
134 6310. .
135 3899. .
136 2473. .
137 1545. .
138 1025. 0
139 711. 0
140 505. 0
141 394. 0
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143 263. 0
144 227. 0
145 193. 0
146 161. 0
147 103. 0
148 130. 0
149 119. 0
150 99. 0
151 111. 0
152 121. 0
153 87. 0
154 101. 0
155 100. 0
156 107. 0
157 107. 0
158 106. 0
159 80. 0
160 100. 0
161 102. 0
162 98. 0
163 82. 0
164 109. 0

167	87.	0
168	84.	0
169	77.	0
170	92.	0
171	71.	0
172	84.	0
173	90.	0
174	75.	0
175	97.	0
176	84.	0
177	77.	0
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191	44.	0
192	66.	0
193	59.	0
194	71.	0
195	40.	0
196	66.	0
197	0.	0
198	0.	0
199	0.	0
200	0.	0
STOP	0	

CHAN SPECT

0
1	16300.
2	16400.
3	16912.
4	16372.
5	16285.
6	15762.
7	16040.
8	16056.
9	15671.
10	15563.
11	15495.
12	15680.
13	15492.
14	15681.
15	15814.
16	15825.
17	16191.
18	16464.
19	17165.
20	19137.
21	20645.
22	21034.
23	19904.
24	18353.
25	17505.
26	17053.
27	17032.
28	17065.
29	16853.
30	17328.
31	17137.
32	16807.
33	16643.
34	16391.
35	16315.
36	16273.
37	16231.
38	16004.
39	15931.
40	15991.
41	16035.
42	16007.
43	16081.
44	16165.
45	16001.
46	16397.
47	16484.
48	16622.
49	16300.
50	16830.
51	16741.
52	17093.

CH-60

55	17027.	*
56	17143.	*
57	17013.	*
58	17443.	*
59	17269.	*
60	17476.	*
61	17605.	*
62	17800.	*
63	18177.	*
64	18277.	*
65	18323.	*
66	18707.	*
67	19033.	*
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70	18293.	*
71	19395.	*
72	19737.	*
73	19922.	*
74	20114.	*
75	20507.	*
76	20675.	*
77	21032.	*
78	20996.	*
79	21140.	*
80	22030.	*
81	21965.	*
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83	23300.	*
84	23386.	*
85	23747.	*
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87	24021.	*
88	24225.	*
89	24213.	*
90	24086.	*
91	23536.	*
92	22753.	*
93	22009.	*
94	20521.	*
95	19252.	*
96	18180.	*
97	17254.	*
98	16303.	*
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127	42603.0
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141	1607.0
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144	591.0
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151	217.0
152	242.0
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157	197.0
158	200.0
159	230.0
160	203.0
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163	200.0
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186	193. 0
187	167. 0
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189	175. 0
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192	191. 0
193	169. 0
194	199. 0
195	165. 0
196	187. 0
197	181. 0
198	187. 0
199	206. 0
200	176. 0
STOP	0

CHAN SPECT

K-42

	CHAN	SPECT	K-42
1	43131.	.	*
2	53006.	.	*
3	49594.	.	*
4	40160.	.	*
5	35607.	.	*
6	31896.	.	*
7	29223.	.	*
8	27060.	.	*
9	24507.	.	*
10	22333.	.	*
11	20355.	.	*
12	19215.	.	*
13	17616.	.	*
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23	10715.	.	*
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25	9385.	.	*
26	8804.	.	*
27	8545.	.	*
28	8555.	.	*
29	9245.	.	*
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108	5293.	.	*

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124	5275.	.	*
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128	3673.	.	*
129	3226.	.	*
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133	2067.	.	
134	1913.	.	
135	1774.	.	
136	1765.	.	
137	1703.	.	
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141	4002.	.	*
142	5989.	.	*
143	8722.	.	*
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146	19193.	.	*
147	21810.	.	*
148	23562.	.	*
149	23859.	.	*
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152	16443.	.	*
153	13103.	.	*
154	9673.	.	*
155	7002.	.	*
156	4800.	.	*
157	3223.	.	*
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161	575.	0	
162	420.	0	
163	296.	0	
164	213.	0	

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200	24.	0
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225 51. 0
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227 36. 0
228 35. 0
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233 42. 0
234 42. 0
235 23. 0
236 32. 0
237 26. 0
238 32. 0
239 45. 0
240 31. 0
241 32. 0
242 26. 0
243 17. 0
244 21. 0
245 26. 0
246 26. 0
247 13. 0
248 13. 0
249 15. 0
250 0. 0

STOP 0

VITA

Jean Paulsen was born in Tampa, Texas, on June 17, 1936. He received his secondary education at Franklin High School, Franklin, Louisiana, in May, 1954, and at the Louisiana State University at Baton Rouge, from February, 1956, until November, 1958. He graduated in December, 1958, with a Bachelor of Science degree in Education. His major field of study was the Behavior of Animals.

He has been employed by the Louisiana State University Board of Education as a teacher, first in the elementary school system and later in the secondary school system.

He has also taught at the University of Southern Mississippi, Hattiesburg, Mississippi, and at the University of Southern Indiana, Evansville, Indiana.

He is currently employed by the University of Southern Indiana, as a teacher in the Department of Curriculum and Instruction.

He is married and has two children, a son, Michael, and a daughter, Linda.

He is a member of the American Association of University Professors.

He is a member of the National Council of Teachers of Mathematics.

He is a member of the National Council of Teachers of Science.

He is a member of the National Council of Teachers of English.

He is a member of the National Council of Teachers of Social Studies.

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He is a member of the National Council of Teachers of Science.

He is a member of the National Council of Teachers of English.

He is a member of the National Council of Teachers of Social Studies.

He is a member of the National Council of Teachers of Mathematics.

EXAMINATION AND THESIS REPORT

VITA

James A. Paulsen
James A. Paulsen

James Paulsen was born in Pampa, Texas, on June 17, 1936. He

completed his secondary education at Franklin High School, Franklin, Louisiana, in May, 1954, and attended Louisiana State University until 1958. He entered the U.S. Navy in 1958 and was discharged in November 1964. He attended Louisiana State University from February, 1965, until May, 1967, at which time he received his Bachelor of Science degree in Chemical Engineering. He entered the Louisiana State University Graduate School in September, 1967, and is at present a candidate for the Master of Science degree in Nuclear Engineering.

James A. Paulsen
James A. Paulsen
James A. Paulsen

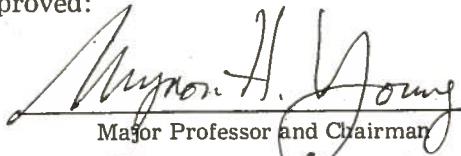
EXAMINATION AND THESIS REPORT

Candidate: James A. Paulsen

Major Field: Nuclear Engineering

Title of Thesis: A Hybrid Computer Filter For Unfolding Gamma-Ray Spectra

Approved:



Myron H. Young

Major Professor and Chairman



Max Goodrich

Dean of the Graduate School

EXAMINING COMMITTEE:



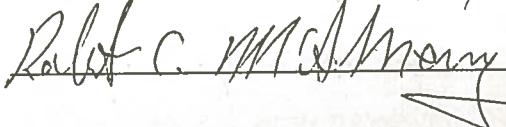
Bert Wilkins Jr.



Frank A. Idsinga

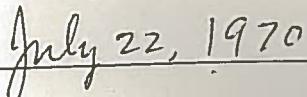


Donald B. Greenberg



Robert C. McMurtry

Date of Examination:



July 22, 1970